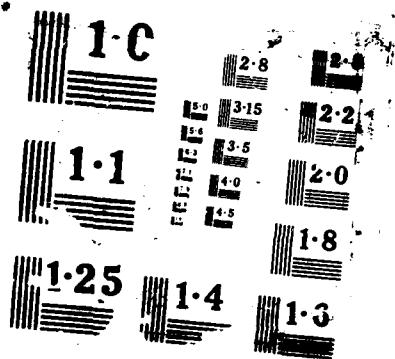


AD-A193 112 WHERE SOLITONS MAY BE HIDING IN DNA AND THEIR POSSIBLE 1/1  
SIGNIFICANCE IN RNA TRANSCRIPTION(U) PURDUE UNIV  
LAFAYETTE IND DEPT OF PHYSICS E W PROHOF SKY 10 MAR 88  
UNCLASSIFIED N00014-86-K-0252 F/G 6/2 NL



2  
DTIC FILE COPY

REPORT DOCUMENTATION PAGE

AD-A193 112

NA

Distribution Unlimited

STANDARD FORM 101-7, 12-75  
1. REPORTING ORGANIZATION NAME AND ADDRESS

Purdue Research Foundation

NA

2. NAME OF FUNDING ORGANIZATION

Purdue Research Foundation

3. ADDRESS (City, State, and ZIP Code)

W. Lafayette, IN 47907

4. OFFICE SYMBOL  
(If applicable)

NA

5. NAME OF FUNDING ORGANIZATION

Technology Applications Programs

6. ADDRESS (City, State, and ZIP Code)

7. ADDRESS (City, State, and ZIP Code)

OSD/SD10/TA  
Washington, D.C. 20301-7100

8. NAME OF FUNDING SPONSORING  
ORGANIZATION Strategic Defense  
Initiative Organization

9. ADDRESS (City, State, and ZIP Code)

Washington, D.C. 20301-7100

10. OFFICE SYMBOL  
(If applicable)

NA

11. FUNDING INSTRUMENT IDENTIFICATION NUMBER

NO0014-86-K-0252

12. SOURCE OF FUNDING NUMBERS

PROJECT ELEMENT NO 63222C	PROJECT NO	TASK NO	ACCESSION NO
---------------------------------	---------------	------------	--------------

13. TITLE (Include Security Classification)

Where Solitons May be Hiding in DNA and Their Possible Significance in RNA Transcription

14. PERSONAL AUTHOR(S):

E.W. Prohofsky

15. TYPE OF REPORT

Technical

16. DATE COVERED

REV 4/81 10 3/88

17. DATE OF REPORT (Year Month Day)

1988 - 3-10

18. PAGE COUNT

19. SUPPLEMENTARY NOTATION

20. COSATI CODES

FIELD	GROUP	SUB GROUP

21. SUBJECT TERMS (Continue on reverse if necessary and identify by block number)

22. ABSTRACT (Continue on reverse if necessary and identify by block number)

We find that the hydrogen bond stretch bands of the double helix appear to be nonlinear enough to support solitary wave energy concentration. Coupling this fact to predictions of our selfconsistent theory of helix melting gives rise to speculations of a mechanism for base pair melting in RNA transcription which is consistent with the known energy needs of that process.

Ribozyme

DTIC  
ELECTED  
APR 11 1988  
S A H

DISTRIBUTION STATEMENT A

Approved for public release;  
Distribution Unlimited

23. DISTRIBUTION AVAILABILITY OF ABSTRACT

UNCLASSIFIED/UNLIMITED  SAME AS RPT  DTIC USERS

24. ABSTRACT SECURITY CLASSIFICATION

U

25. NAME OF RESPONSIBLE INDIVIDUAL  
Charles L. Houston III

26. TELEPHONE NUMBER (Area Code) (202) 693-1556  
27. OFFICE SYMBOL

OSD/SD10/TA

Where Solitons May Be Hiding in DNA and Their Possible Significance  
in RNA Transcription

E. W. Prohofsky

Department of Physics  
Purdue University  
West Lafayette, IN 47907

We find that the hydrogen bond stretch bands of the double helix appear to be nonlinear enough to support solitary wave energy concentration. Coupling this fact to predictions of our selfconsistent theory of helix melting gives rise to speculations of a mechanism for base pair melting in RNA transcription which is consistent with the known energy needs of that process.

Several recent papers have discussed the possibility of solitary waves in double helical DNA associated with the helical acoustic modes (1). The work was stimulated by observations in DNA of resonant microwave absorption (2). Papers have also appeared which have disputed the theoretical predictions of the likelihood of such solitary waves (3) and papers have also appeared which challenge the experimental observations (4). Solitary wave phenomena require nonlinear effects and both experimental observations (5) and theoretical analysis (6) indicate that helix acoustic modes are quite linear in their dispersion and amplitude dependence. In our theoretical studies of helix lattice dynamics of DNA we have found that a different region of the DNA spectrum is far more nonlinear (7) and seems to satisfy the quantitative requirements for solitary wave effects.

The highly nonlinear modes are the hydrogen bond stretch modes at 16 - 120  $\text{cm}^{-1}$ . These modes have been observed in both low frequency Raman scattering (8) and Fourier transform infrared absorption (9) experiments. In both cases the modes have been seen to soften in frequency with increasing temperature. We have predicted theoretically that these modes are temperature sensitive in agreement with observations (10).

The same vibrational modes have played a central role in the theory of helix melting that we have developed based on helical lattice dynamics. The temperature dependence discussed above, coupled with thermal expansion, also lead naturally to hydrogen bond breakup and helix melting. These results for poly(dG)•poly(dC) predict melting at a temperature in fair agreement with observations. In that work all the parameters were fitted to spectral data at room temperature and no parameters were fitted to melting. (11)

Distribution/	
Availability Codes	
Dist	Avail and/or Special
A-1	

Helix melting is an essential element in many biological processes. For example, it is essential in the process of RNA transcription. Recent reviews of this process reaffirm the belief that energy must flow into the RNA polymerase complex from its surroundings (12). The ability of solitary wave effects to concentrate excitations in these bands would help in understanding the problems with energy flow in these processes. This biological role for nonlinear effects is in line with the original suggestions of solitary wave significance in biological processes (13).

In this paper we discuss our helical lattice dynamics results for the nonlinear modes and show how these results can be used to extract parameters needed to evaluate the possibilities of nonlinear concentration. We then use the theory of melting developed to give order of magnitude estimates of the amplitudes needed to induce base pair melting, and discuss the significance of these results for RNA transcription.

#### 3. Hydrogen Bond Stretch Modes in the Double Helix

In our study of the vibrational modes of the double helix we have found several bands which can be characterized as vibrations that compress the H-bonds between base pairs. There are overlapping bands for all combinations of base pairs and several pass bands in this region should exist regardless of base sequence (14). Dispersion curves, i.e., the change in frequency with wave vector are shown for several such bands in Figure 1 (15). Theta is the angle between the displacement from one unit cell to the next. These curves have local minimum frequencies. In one case the local minimum is at the middle of a line in other cases its at a zone edge. The curves could

be extended about either end i.e. the total  $\theta$  range can run from  $-180^\circ < \theta \leq 180^\circ$  or  $0^\circ \leq \theta < 360^\circ$ . The curves would have symmetry about the extension as the positive and negative phase shifts are connected by time reversal invariance. Correlating frequency with energy, and phase shift with wavevector and hence momentum for these excitations, one can relate these curves to particles with energy vs. momentum relations approximated by

$$E = E_0 + \frac{(\Delta p)^2}{2m} \quad (1)$$

The effective mass  $m$  is then

$$m = \frac{(\Delta p)^2}{2\Delta E}$$

$$m = \frac{(\hbar^2 \Delta q^2)}{2\hbar \omega_d}$$

$$\frac{m}{\hbar} = \frac{\Delta q^2}{2\Delta \omega_d}$$

if  $a$  is the lattice spacing

$$\frac{ma^2}{\hbar} = \frac{(\Delta \theta)^2}{2\Delta \omega_d} \quad (2)$$

where  $\theta$  is in radians and  $\Delta \omega_d$  is the change in frequency due to dispersion effects for a shift in the phase from cell to cell of  $\Delta$ . In the solitary wave analysis by Davydov (13) the wave equation for such a system is given by the nonlinear Schrödinger Equation in one dimension

$$(i\hbar \frac{\partial}{\partial t} + \frac{\hbar^2}{2m} \frac{\partial^2}{\partial z^2} + G|\psi|^2) \psi(z, t) = 0 \quad (3)$$

with the substitutions

$$x = z/a \quad \tau = \frac{\hbar t}{2ma^2} \quad \text{and} \quad g = \frac{ma^2 c}{\hbar^2} \quad (4)$$

this becomes

$$(i \frac{\partial}{\partial \tau} + \frac{\partial^2}{\partial x^2} + 2g|\psi|^2) \psi(x, \tau) = 0 \quad (5)$$

which has the solitary wave solution

$$\psi = \frac{1}{2} \sqrt{g} e^{i(2kx - \omega \tau)} \operatorname{sech} \left[ \frac{g(x - 4k\tau)}{2} \right] \quad (6)$$

with

$$\omega = 4(k^2 - \frac{g^2}{16}) \quad k = \frac{maV}{\hbar} \quad (7)$$

where  $V$  is the velocity of the centroid of the solitary wave. Concentration of the excitation occurs in the square wave excitation case (13) for  $g > g_{cr}$  where

$$g_{cr} \approx \frac{\pi^2}{4k} \quad (8)$$

and  $\tau$  is the extent of the initial excitation along  $x$ .

The term  $2g|\psi|^2$  is a measure of how the energy is lowered by increased concentration of the probability amplitude. In our selfconsistent temperature dependent lattice calculations we have found that the hydrogen bond stretch modes are anharmonic, i.e., they lower in frequency as the excitation level or temperature increase, we have found this shift in frequency to be in rough agreement with the observed frequency shift from low temperature to room temperature (10). This decrease in energy with increase in concentration of amplitude provides the selftrapping effect which enhances concentration of the excitation. In figure 2 we show a calculation (14) of the frequency of several modes in poly(dL)-poly(C) as a function of temperature for temperatures above

room temperature. For vibrational systems the mean wave amplitude squared is proportional to the temperature. The frequency is seen to drop with increasing temperature and hence mean squared amplitude.

We can estimate  $G$  for particular excitation modes from Figure 2. In Equation 3,  $|\psi|^2$  is a normalized nondimensional squared wave amplitude. We associate  $\Delta |\psi|^2$  with  $\frac{\Delta T}{T}$  for regions of curve in Figure 2. We also associate the appropriate  $\Delta E$  with  $\hbar\Delta\omega_A$  over that  $\Delta T$ .  $\Delta\omega_A$  is the change in frequency due to anharmonicity not dispersion. Then:

$$\frac{G}{\hbar} = \frac{\Delta\omega_A}{\Delta T/T} \quad (9)$$

Using this and Eqs. (2) in Eqn. (4)

$$g = \frac{ma^2}{\hbar} \times \frac{G}{\hbar} = \frac{(\Delta\theta)_d^2}{2\Delta\omega_d} \frac{\Delta\omega_A}{\Delta T/T} \quad (10)$$

For the  $85 \text{ cm}^{-1}$  band from Figures 1 and 2 at  $T \approx$  room temperature we estimate  $a \approx 10$ . For reasons discussed in the next section we estimate  $l_{\min}$  the minimum initial size of the excitation to be 2.5 in units of  $a$ . For this value  $g_{cr} \approx 1$ . One can expect nonlinear concentration of excitations in these bands in DNA. In Figure 2 one can see that  $G$  would be even larger at the higher levels of excitation needed for helix melting.

### III. Discussion

From the review of RNA transcription the basic process, a "step", fuses an RNA nucleotide to the growing RNA transcript. For each step a DNA base pair is split on the downstream side, a DNA + RNA hybrid base pair is formed, a DNA + RNA hybrid base pair is split, a DNA base pair upstream is refolded, and

a polymerization event in the RNA backbone takes place. As a result of the backbone chemistry a pyrophosphate group is removed from the nucleoside triphosphate (NTP) which is the source of the added RNA section. The pyrophosphate group is later hydrolyzed but not at a site directly connected to the other operations. The overall energy flow in the transcription complex is such that on the average 0.5 kcal/mole transcribed must be absorbed by the complex from its surroundings.

For the purposes of this discussion we will explore ways in which energy introduced into the complex can bring about base pair melting. The free energy needed to melt a GC pair is generally accepted to be 3.5 kcal/mole and that for an AT pair 1 kcal/mole (12). If inflow of this amount of energy occurred the net energy requirements of transcription would easily be met. The reason to consider this form of energy transfer to the transcription complex is that we believe it would involve the nonlinear hydrogen bond stretch modes. In our theory of melting we calculate the vibrational fluctuations  $D$  on each hydrogen bond where  $D = \langle u^2 \rangle - \langle u \rangle^2$  and  $u$  a is the displacements of atoms from equilibrium due to phonon amplitude. The selfconsistent softening of the bond and thermal expansion are calculated based on  $D$ . The softening and expansion lead to a melting of the hydrogen bonds. In this sense  $D$ , which grows with excitation level, drives the melting. In our model where we looked at melting adjacent to a melted base pair in poly(dC)•poly(dC) we found that when a hydrogen bond has  $D = 0.040 \text{ \AA}^2$  melting occurred (11).

We also investigated the contribution of individual vibrational bonds  $D(u)$  to the total fluctuation  $D$  as seen in Figure 3. The hydrogen bond stretch bond

at  $95 - 100 \text{ cm}^{-1}$  at 340K has  $D(\omega) \approx 5.5 \times 10^{-3} \text{ \AA}^2$  from the entire band. Ignoring the contributions of other bands one would only have to increase the excitation of this band by a factor of 7.27 to bring on melting from the energy in this band alone. Since the thermal occupation is about 2 phonons at this temperature one would require  $\approx 14-15$  phonons to bring about melting. The energy of  $\approx 14$  phonons is approximately the 3.5 kcal/mole generally associated with the binding energy of GC base pairs. The localization to the hydrogen bond of these excitations is in the localized amplitudes calculated and plotted in Figure 3.

If we consider melting a GC base pair in terms of the needed energy in one mode while all others are in thermal equilibrium at room temperature a different estimation is appropriate. One still needs a total  $D \approx 0.004 \text{ \AA}^2$ . The thermal activation of all modes at room temperature is  $D_R \approx 0.028 \text{ \AA}^2$  (7). The needed additional fluctuation is  $\Delta D = 0.012 \text{ \AA}^2$ . For  $D(\omega)$  of the  $95-100 \text{ cm}^{-1}$  band  $D(\omega) = 0.0055$  to make up this difference the excitation level would have to be 2.8 times that at 343K. This requires  $\approx 6$  phonons in a packet arriving at the base pair adjacent to the melted region. The thermal activation level is  $\approx 2$ . The mean large fluctuation value from Bose-Einstein statistics is  $\approx 4-5$ . Such a concentration of energy may be achieved by capturing energy from the surroundings as is apparently needed to explain the energetics of DNA transcription. Such an excitation could travel in unmelted DNA without disrupting the DNA as the  $D(\omega)$  for all bands is less for base pairs far from a melted section. However the  $95-100 \text{ cm}^{-1}$  band is about half the value at a cell(+) as for cell (+1) next to the melt, as seen in Figure 3. When the energy propagates into the cell adjacent to the melt the change in localized mode behavior greatly increases fluctuation amplitudes and disrupts the base pair.

It is an interesting question as to whether the RNA polymerase is a proper thermodynamic engine which runs on a directed flow of energy or is driven by thermal fluctuations as many believe (12). If solitary wave effects can be shown to in fact exist in DNA one will have come a long way to showing that the RNA polymerase can be a proper thermodynamic engine. Energy from the hydrolysis of pyrophosphate could be absorbed before the thermalization of this energy. The entire helix could act as an antenna for the absorption greatly increasing the fraction of energy absorbed before thermalization. The processing and concentration of energy would occur as the energy propagated on the one dimensional helix to the region of transcription.

In the earlier section we estimated  $\ell$  the initial size of the excitation as being at least  $\approx 2.5$  base units in length. This is based on the idea of a planar longitudinal wave or shock wave from pyrophosphate hydrolysis impinging on a section of DNA helix from the side. A helical structure is just what's needed to convert such an external excitation into a short wavelength excitation in the internal coordinates of the DNA because of  $k$ -conservation selection rules. For some forms of excitation a 10 fold helix would select excitations of 10 base pair wavelength or  $\theta = 36^\circ$  on dispersion curves such as Figure 1. To generate a hydrogen bond stretch wave where compression of either strand leads to compression of the hydrogen bonds the wavelength would be 5 base pairs or  $72^\circ$ . We have associated half this wavelength with the minimum extent of the initial square wave excitation and set  $\ell = 2.5$ . Short wavelengths can lead to more localized wave packets and localization or concentration is needed to efficiently melt base pairs adjacent the already melted area. It is interesting

to speculate on how the existence of a small integer-fold helical structure for DNA interacts with the need for energy transfer in process such as transcription in early evolution of the system.

Better nonlinear models of wave propagation more appropriate to the case of dispersive optical modes would allow a more quantitative estimate of many of the energies which have been crudely estimated in this phenomenological discussion. The role of DNA excitations in biological processes may be important and seems likely to be, based on these simple estimates.

This work supported in part by ONR Grant N00014-86-K-0252 and NIH Grant GM 24443.

## References

1. A. C. Scott and J. H. Jensen, *Phys. Lett.* 109A 243-245 (1985).
2. G. S. Edwards, C. C. Davis, J. D. Saffer, and M. L. Swicord, *Phys. Rev. Lett.* 53, 1284 (1984).
3. L. L. Van Zandt, *Private Communication*.
4. C. Gabriel, E. H. Grant, R. Tata, P. R. Brown, B. Gestblom and E. Noreland, *Nature* 328, 145-146 (1987).
5. S. M. Lindsay and J. Powell, *Structure and Dynamics - Nucleic Acids and Proteins*, E. Clementi and R. H. Sarma Eds. Adenine Press, New York, pp. 241-260.
6. W. N. Mei, M. Kohli, E. W. Prohofsky and L. L. Van Zandt, *Biopolymers* 20, 833-852, (1981).
7. Y. Kim, K. V. Devi-Prasad, and E. W. Prohofsky, *Phys. Rev. B* 32, 5185-5189 (1985).
8. F. Urabe and Y. Tominaga, *J. Phys. Soc. Japan*, 50, 3543-3544 (1981).
9. J. W. Powell, G. S. Edwards, L. Genzel, and A. Wittlin, *Phys. Rev. A* 35, 3929 (1987).
10. Y. Kim and E. W. Prohofsky, *Phys. Rev. B* 36, 3449 (1987).
11. Y. Kim and E. W. Prohofsky, *Phys. Rev. B* 33, 2399 (1987).
12. T. D. Yager and P. H. von Hippel to appear in *Escherichia coli and Salmonella typhimurium: cellular and Molecular Biology*, F. Neidhart et al. eds. American Society of Microbiology.
13. A. S. Dovydov, *Solitons in Molecular Systems*, D. Reidel Pub. Co. Dordrecht (1985).

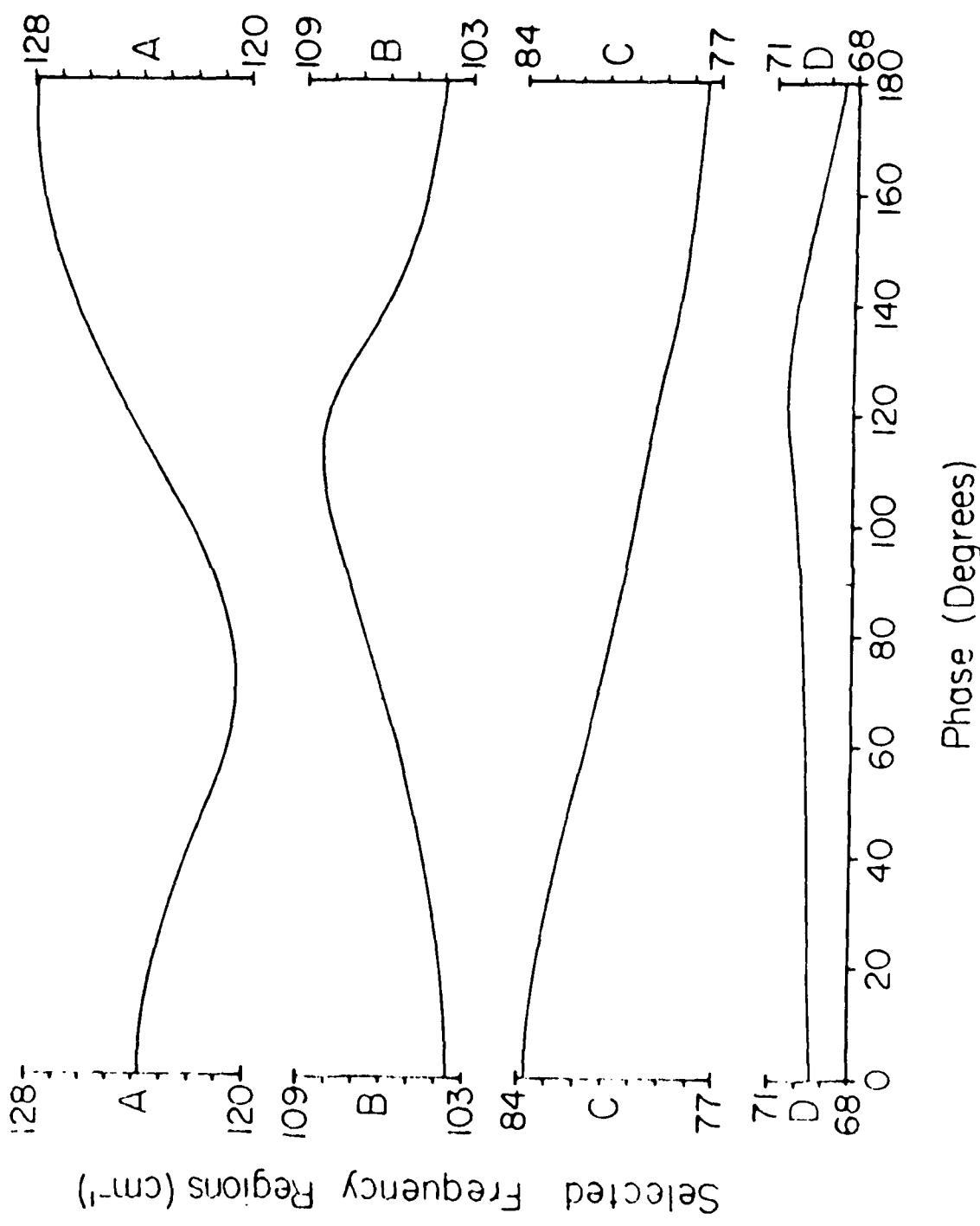
14. Y. Gao, K. V. Devi-Prasad and E. W. Prohofsky, *J. Chem. Phys.* 80, 6291 (1984).
15. E. W. Prohofsky, *Comments Mol. Cell. Biophys.* 2, 65 (1983).

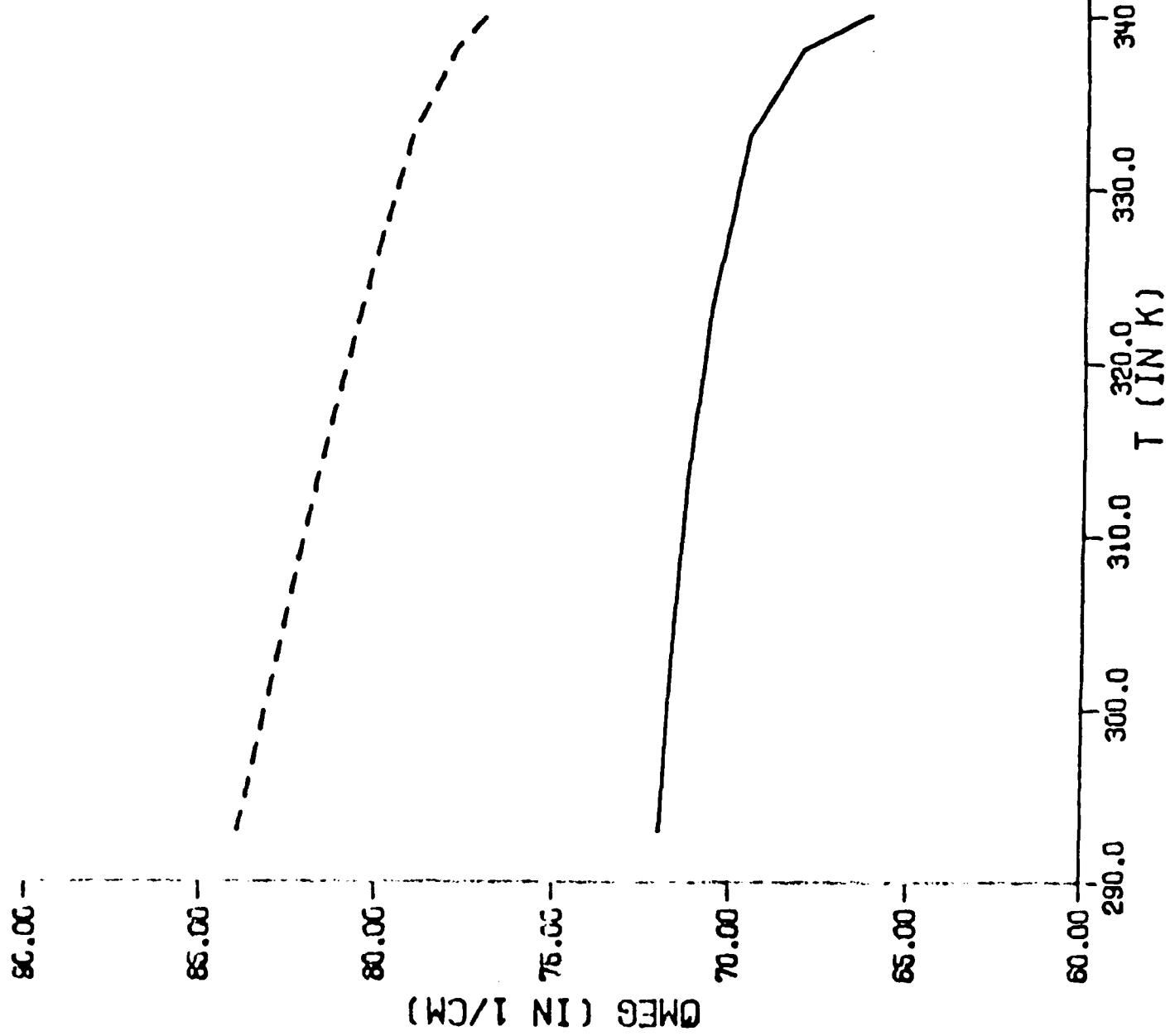
## Figure Captions

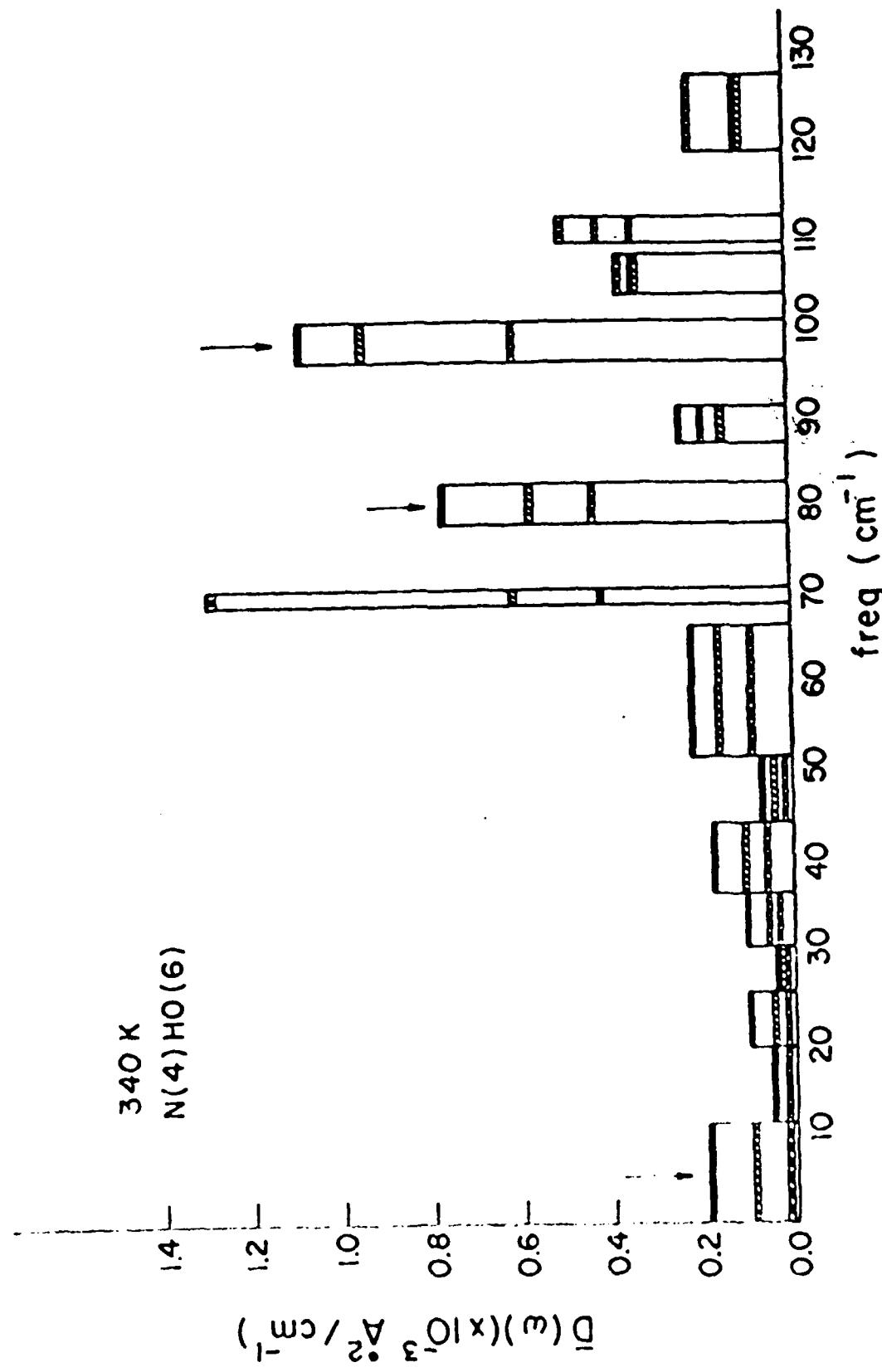
Figure 1. Dispersion relations for several optical bands in poly(dG)·poly(dC). The frequency is in  $\text{cm}^{-1}$  and the phase is the relative phase between neighbor unit cells. In terms of a wavevector  $q$ , this phase is  $qa$  where  $a$  is the lattice constant.

Figure 2. Calculated change in frequency of selected zone center modes of poly(dG)·poly(dC) as a function of temperature. For details see reference (14).

Figure 3. The contribution to the fluctuation amplitude by individual vibrational bands for the hydrogen bond adjacent the major groove in poly(dG)·poly(dC). The three divisions are for the two base pairs adjoining a melted base pair on either side and for a base pair far away from the melted base pair. For the  $95-100 \text{ cm}^{-1}$  band the highest division is for the (+1) base pair i.e.  $3' \rightarrow 5'$  direction in the G chain. The next highest is for the (-1) base pair i.e. on the other side of the melt. The lowest division is for the ( $\infty$ ) base pair i.e. far away from the melt. See reference 11 for details.







END  
DATE  
FILMED  
DTIC

JULY 88